# Lecture 18: Finite Element Methods (1)

Today:

- *Finite element methods* for *boundary value problems* (*BVPs*)
  - This is a spectral method too!
  - But it is based on *locally* defined functions, not global ones

# Where are we up to now?

Last time.

- (A) We developed a framework for approximating the solution to a BVP using global spectral methods NOTE: starting index at 1, not 0; ending
- (B) We arrived at a linear system to solve for the coefficients

NOTE: starting index at 1, not 0; ending index at *n* + 1, not *n*. Either approach is ok as long as you are consistent!

$$\Rightarrow \begin{bmatrix} (b_{1}, b_{1})_{E} & (b_{2}, b_{1})_{E} & \cdots & (b_{n}, b_{1})_{E} & (b_{n+1}, b_{1})_{E} \\ (b_{1}, b_{2})_{E} & (b_{2}, b_{2})_{E} & \cdots & (b_{n}, b_{2})_{E} & (b_{n+1}, b_{2})_{E} \\ \vdots & \vdots & \cdots & \vdots & \vdots \\ (b_{1}, b_{n})_{E} & (b_{2}, b_{n})_{E} & \cdots & (b_{n}, b_{n})_{E} & (b_{n+1}, b_{n})_{E} \\ (b_{1}, b_{n+1})_{E} & (b_{2}, b_{n+1})_{E} & \cdots & (b_{n}, b_{n+1})_{E} & (b_{n+1}, b_{n+1})_{E} \end{bmatrix} \begin{bmatrix} c_{1} \\ c_{2} \\ \vdots \\ c_{n} \\ c_{n+1} \end{bmatrix} = -\begin{bmatrix} (f, b_{1})_{s} \\ (f, b_{2})_{s} \\ \vdots \\ (f, b_{n})_{s} \\ (f, b_{n+1})_{s} \end{bmatrix}$$
(1)  
Call this **Gc** = **b**

(C) Can back out our approximate solution via

$$u_a(x) = \sum_{j=1}^{n+1} c_j b_j(x)$$

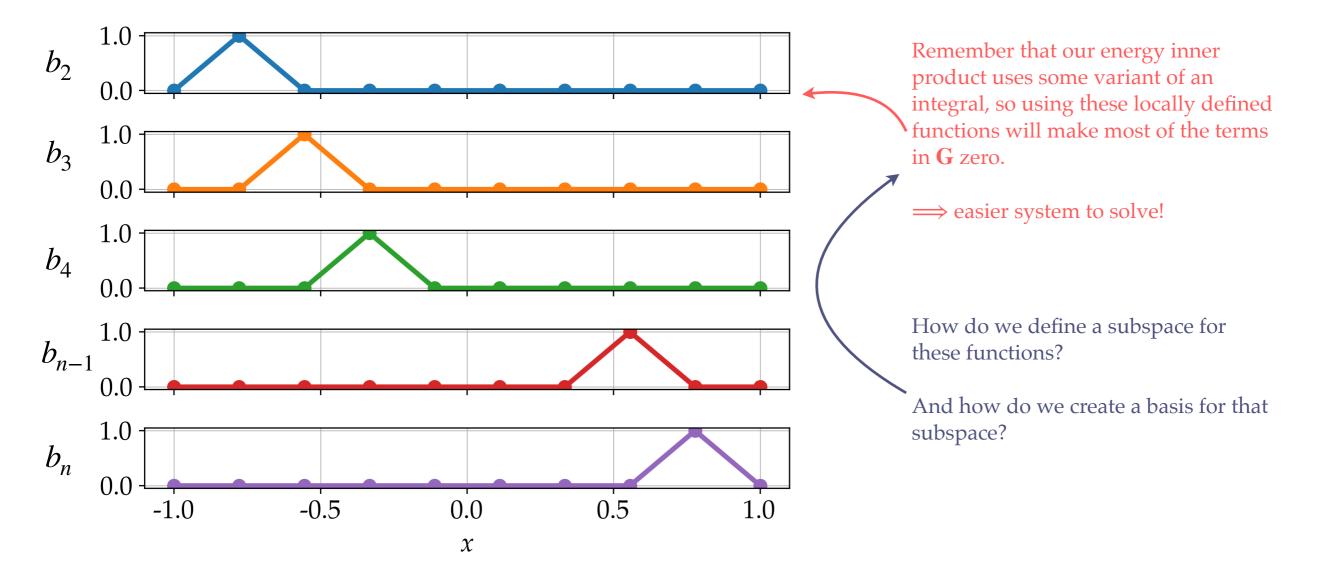
*Today.* The *finite element method (FEM):* a locally-based spectral method

- What is the premise for the finite element method?
- How do we define a subspace that admits locally defined functions?
- How do we pick a basis for that subspace, and how does that affect (1)

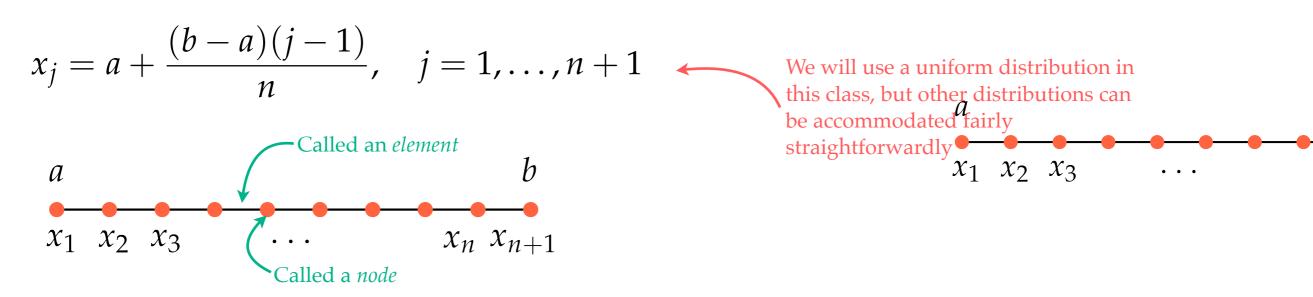
## The premise of FEMs

Philosophy behind the finite element method

The finite element method is a spectral method that uses locally defined functions to create a matrix *G* with predominately zero entries to facilitate a fast solution of Gc = b.



#### Define the subspace using a discretized domain



Now that we have a set of points, we can define a function that is piecewise-linear using these points:

$$g(x) = a_i x + d_i$$
  $x \in [x_{i-1}, x_i], i = 2, ..., n+1$ 

This gives us intuition for the subspace we want! Let's define our subspace to be the collection of functions that

(A) Satisfy the BCs

(B) Are piecewise-linear on the set of discrete points:

$$\mathcal{V}_n^L = \{g(x) : g(a) = g(b) = 0 \text{ and} \\ g(x) = a_i x + d_i, \text{ for } a_i, d_i \in \mathbb{R}, x \in [x_{i-1}, x_i], i = 2, \dots, n+1\}$$
  
Brain teaser — how would you go about showing this is a subspace?!

A basis for 
$$\mathcal{V}_n^L$$

We will show that the functions pictured in slide 3 form a basis for  $\mathcal{V}_{n}^{L}$ , define them mathematically, and then use them to build (1) for our FEM

The basis functions for  $\mathcal{V}_n^L$  are defined as

Notice that  

$$b_{i}(x_{j}) = \begin{cases} 1 & i = j \\ 0 & i \neq j \end{cases} \quad b_{i}(x) = \begin{cases} \frac{1}{\Delta x} [x - a - (i - 2)\Delta x], & \text{if } x \in [x_{i-1}, x_{i}] \\ -\frac{1}{\Delta x} [x - a - i\Delta x], & \text{if } x \in [x_{i}, x_{i+1}] \end{cases} \quad i = 2, \dots, n$$

$$0, \qquad \text{else}$$
We exclude  $i = 1, n + 1$ , because we will show

 $\mathscr{B} = \{b_2, ..., b_n\}$  forms a basis for  $\mathscr{V}_n^L$ :

We exclude i = 1, n + 1, because we will show these are taken care of by the BCs

The functions are clearly linearly independent (see the fig on slide 3 — can't write any one of the functions as a linear combo of the others!)

Take any piecewise linear function q(x) defined as a set of lines over  $[x_1, x_2], [x_2, x_3], \dots, [x_n, x_{n+1}]$ . Then this q(x) can be written in terms of the basis functions and its *nodal values*  $q(x_1), q(x_2), \dots, q(x_n)$ 

$$q(x) = \sum_{j=2}^{n} q(x_j) b_j(x)$$

OK, so we have our basis functions! How does the matrix system (1) simplify for this basis? Let's figure that out for the 1D Poisson problem with zero-Dirichlet BCs!

### Simplifying (1) for our FEM basis

Recall that our energy inner product is defined in terms of an integral of derivatives. So let's look at the derivative of the basis functions:

$$b_i'(x) = \begin{cases} \frac{1}{\Delta x} & x \in [x_{i-1}, x_i] \\ -\frac{1}{\Delta x} & x \in [x_i, x_{i+1}] \\ 0 & else \end{cases}$$

Because these derivatives are only nonzero for a small sub-interval of the domain, the only terms of row *i* of **G** that survive are  $(b_i, b_{i-1})_E$ ,  $(b_i, b_i)_E$ ,  $(b_i, b_{i+1})_E$ 

These terms can be computed analytically:

$$(b_i, b_{i-1})_E = \int_{a+(i-1)\Delta x}^{a+i\Delta x} \left(\frac{-1}{\Delta x}\right) \left(\frac{1}{\Delta x}\right) dx = -\frac{1}{\Delta x}$$

$$(b_i, b_i)_E = \int_{a+(i-1)\Delta x}^{a+(i+1)\Delta x} \left(\frac{1}{\Delta x}\right) \left(\frac{1}{\Delta x}\right) dx = \frac{2}{\Delta x}$$

$$(b_i, b_{i+1})_E = \int_{a+i\Delta x}^{a+(i+1)\Delta x} \left(\frac{-1}{\Delta x}\right) \left(\frac{1}{\Delta x}\right) dx = -\frac{1}{\Delta x}$$

Similarly, the  $(f, b_i)_s$  term reduces to

$$(f, b_i)_s = \int_a^b f(x)b_i(x)dx = \int_{a+(i-1)\Delta x}^{a+(i+1)\Delta x} f(x)b_i(x)dx$$

### Rewriting (1) for our FEM basis

$$\frac{1}{\Delta x} \begin{bmatrix} -2 & 1 & \cdots & 0 & 0 \\ 1 & -2 & \cdots & 0 & 0 \\ \vdots & \vdots & \cdots & \vdots & \vdots \\ 0 & 0 & \cdots & -2 & 1 \\ 0 & 0 & \cdots & 1 & -2 \end{bmatrix} \begin{bmatrix} u_2 \\ u_3 \\ \vdots \\ u_{n-1} \\ u_n \end{bmatrix} = \begin{bmatrix} \int_a^{a+2\Delta x} f(x)b_2(x)dx \\ \vdots \\ \int_{a+\Delta x}^{a+3\Delta x} f(x)b_3(x)dx \\ \vdots \\ \int_{a+(n-2)\Delta x}^{a+(n)\Delta x} f(x)b_{n-1}(x)dx \\ \int_{a+(n-1)\Delta x}^{a+(n+1)\Delta x} f(x)b_n(x)dx \end{bmatrix}$$

*Punchline:* lots of zeros in that matrix; easier to solve than the generic form of (1)

*Punchline:* Solve matrix system  $\implies$  get the  $u_1, u_2, ..., u_n$  that approximate  $u(x_1), u(x_2), ..., u(x_n)$ 

*Next time:* some coding examples and concept discussions!